Toward an efficient hybrid method for pricing barrier options on assets with stochastic volatility

Alexander Lipton^{*}, Artur Sepp[†]

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Abstract

We combine the one-dimensional Monte Carlo simulation and the semianalytical one-dimensional heat potential method to design an efficient technique for pricing barrier options on assets with correlated stochastic volatility. Our approach to barrier options valuation utilizes two loops. First we run the outer loop by generating volatility paths via the Monte Carlo method. Second, we condition the price dynamics on a given volatility path and apply the method of heat potentials to solve the conditional problem in closed-form in the inner loop. We illustrate the accuracy and efficacy of our semi-analytical approach by comparing it with the twodimensional Monte Carlo simulation and a hybrid method, which combines the finite-difference technique for the inner loop and the Monte Carlo simulation for the outer loop. We apply our method for computation of state probabilities (Green function), survival probabilities, and values of call options with barriers. Our approach provides better accuracy and is orders of magnitude faster than the existing methods. s a by-product of our analysis, we generalize Willard's (1997) conditioning formula for valuation of path-independent options to path-dependent options and derive a novel expression for the joint probability density for the value of drifted Brownian motion and its running minimum.

Keywords: barrier options; stochastic volatility; Heston model; heat potentials; semi-analytical solution; Volterra equation; Willards's formula;
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1 Introduction

By expressing prices of European calls and puts in terms of the price of the underlying asset and its volatility, Black and Scholes (1973) and Merton (1973) started the quantitative finance revolution. Their formula, which is known as the Black-Scholes-Merton formula, is based on two assumptions: (a) the risky

^{*}Abu Dhabi Investment Authority, 211 Corniche, PO Box 3600, Abu Dhabi, United Arab Emirates

[†]Sygnum Bank, Uetlibergstrasse 134a, 8045 Zürich, Switzerland

price dynamics can be delta-hedged so that options can be valued using the risk-neutral measure under which the underlying grows at the risk-neutral rate; (b) price evolution is driven by a geometric Brownian motion of the form:

$$\frac{dS_t}{S_t} = rdt + \sigma dW_t, \quad S_0 = s, \tag{1}$$

where W_t is a Brownian motion, r and σ are constant risk-neutral interest rate and volatility of returns, respectively. Thus, at time T, the price S_T has the log-normal distribution:

$$S_T = e^{\left(r - \frac{1}{2}\sigma^2\right)T + \sigma W_T} S_0, \tag{2}$$

Under these assumptions, it is very easy to derive the price of, say, a call option with maturity T and strike K, with payoff of the form $(S_T - K)_+$:

$$C^{BS}(0, S_0, T, K; r, \sigma) = S_0 \mathfrak{N}(d_+) - e^{-rT} K \mathfrak{N}(d_-), d_{\pm} = \frac{-\ln(K/S_0) + rT \pm \sigma^2 T/2}{\sigma \sqrt{T}},$$
(3)

where $\mathfrak{N}(x)$ is the cdf for the standard (0, 1) normal random variable. Lipton (2002) showed that in many situations it is more convenient to write C^{BS} as a Fourier integral:

$$C^{BS}(0, S_0, T, K; r, \sigma) = S_0 \left(1 - \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{e^{(i\chi + 1/2)(\ln(K/S_0) - rT) - (\chi^2 + 1/4)\sigma^2 T/2}}{(\chi^2 + 1/4)} d\chi \right).$$
(4)

However, it became clear in a few years that the Black-Scholes formula, taken literally, is not as helpful as initially thought since it could not reproduce option market prices by using constant volatility for pricing options with different strikes and maturities. The volatility skew (or smile) effect (i.e., the need to use different volatilities to reproduce market prices of European calls and puts with different strikes and maturities) is observed in all markets. Therefore, it is of great interest to practitioners and academics alike. Over the thirty years, many models were developed to describe the smile effect.

Eventually, two approaches emerged - a reduced approach replacing the constant volatility, central to the Black-Scholes approach, by the maturity- and strike-dependent implied volatility, used to match the corresponding market prices. Specifically, instead of a constant σ for all T, K in Eq. (3), a function $\sigma(T, K)$ is used. The corresponding call price has the form

$$C(0, S_0, T, K) = C^{BS}(0, S_0, T, K; r, \sigma(T, K)).$$
(5)

While useful in some situations, this approach is conceptually fairly limited, since it does not explain why volatility is not constant. Hence, a structural approach replacing the geometric Brownian motion with a different, but still risk-neutral, driver for the underlying was quickly developed. For instance, Derman and Kani (1994), Dupire (1994), and Rubinstein (1994) simultaneously and independently developed the local volatility model. Merton (1976), Andersen and Andreasen (2000), Lewis (2001) and many others developed jump-diffusion models. Stochastic volatility models were developed by Hull and White (1987), Scott (1987), Wiggins (1987), Stein and Stein (1991), Heston (1993), Lewis (2000), Bergomi (2015), and others. Various combinations of the above were proposed, for example by Dupire (1996), Jex *et al.* (1999), Hagan *et al.* (2002), Lipton (2002), which culminated in Lipton's universal volatility model; see Lipton (2002). The above-mentioned models are compared and contrasted in Lipton (2002).

Lipton and McGhee (2002) explained that the actual worth of a structural model is not in its ability to price vanilla options, which all structural models worth their salt can do well, but in producing consistent prices for *both* vanillas and first-generation exotics. Despite more than thirty years of strenuous efforts, finding a proper theoretical framework and implementing it in practice remains a significant challenge.

Pricing of exotic options in the presence of a smile is usually difficult and seldom can be done analytically. Asymptotic methods developed by Hull and White (1987), Hagan et al. (2002), Lipton (1997), Lipton (2001), among others, proved to be very useful for solving the corresponding pricing problems. Numerical methods, such as the Monte Carlo simulation (MCS) and the finite difference method (FDM), are equally useful: see Glasserman (2004), Achdou and Pironneau (2005). Adding new methods to the classical ones is definitely worth the effort. One can reduce many problems we wish to solve to the initialboundary value problems (IBVPs) for one-dimensional parabolic partial differential equations (PDEs) with moving boundaries and (or) time-dependent coefficients. Such problems appear naturally in various areas of science and technology. Finding their semi-analytical solutions requires using sophisticated tools, such as the method of heat potentials (MHP) and a complementary method of generalized integral transforms (MGIT). Such methods were actively developed by the Russian mathematical school in the 20th century; see Kartashov (2001) and references therein.

In the context of mathematical finance, A. Lipton and his coauthors actively utilized the MHP; see Lipton (2001), Lipton *et al.* (2019) Lipton and Kaushansky (2020) and references therein. In addition, A. Itkin and his coauthors used the MGIT to price barrier and American options in the semi-closed form; see, e.g., Carr *et al.* (2020), Itkin and Muravey (2021), Carr *et al.* (2022), Itkin and Muravey (2022). In principle, the MHP and MGIT can be generalized and used for any linear differential operator with time-independent coefficients.

The MHP and MGIT boil down to solving linear Volterra equations of the second kind and representing option prices as one-dimensional integrals. Itkin *et al.* (2021) described the MHP and MGIT in the recent comprehensive book and showed that they are much more efficient and provide better accuracy and stability than the existing methods, such as the backward and forward FDM or MCS.

This paper revisits the classical problem of pricing barrier options on assets

with stochastic volatility. We show that by using the concept of conditional independence, we can reduce it to solving an initial-boundary value problem with time-dependent coefficients and subsequent averaging over the space of variance trajectories. Based on this observation, we develop an efficient method that combines the MHP and MCS and provides a fast and accurate solution to the problem at hand. Our method is very general and can handle all known stochastic volatility models, as well as models with rough volatility.¹

The paper is organized as follows. In Section 2, we introduce generic stochastic volatility models and describe the splitting method, which allows one to study the dynamics of the log-price $X_t = \ln (S_t/S_0)$, as a conditionally-independent one-dimensional processes. We specialize these equations for the Heston and Stein-Stein models. We also present the exact Lewis-Lipton and conditional Willard formulas for vanilla options, such as European calls and puts on assets with stochastic volatility and compare the corresponding prices. In Section 3, we introduce barrier options on assets with stochastic volatility, which are the main object of our study. We derive a conditional valuation formula for such options, which generalizes the Willard formula for vanilla options. In Section 4, we describe a hybrid method for pricing barrier options, which relies on the conditional independence decomposition. The method consists of the outer Monte Carlo loop and the inner loop, which requires solving the advection-diffusion problem for the drifted Brownian motion with time-dependent coefficients on a semi-axis. We solve the latter problem via two complementary methods: the FDM and the MHP. Results produced by both methods are in perfect agreement. However, as expected, the second method is orders of magnitude faster than the first one. In Section 4, we use the MHP to solve an old problem in probability theory and show how to find the joint probability density for the value of drifted Brownian motion and its running minimum via the MHP. We draw our conclusions in Section 6.

2 Stochastic volatility models

2.1 Conditionally independent dynamics

We consider the joint evolution of the asset price, S_t , and its stochastic variance, V_t , as follows:

$$\frac{dS_t}{S_t} = rdt + \sqrt{V_t} \left(\rho dB_t + \sqrt{1 - \rho^2} dW_t \right), \quad S_0 = s,
dV_t = \Phi \left(V_t \right) dt + \Psi \left(V_t \right) dB_t, \quad V_0 = v,$$
(6)

where B_t and W_t are independent Brownian motions. Given that Eqs (6) are scale invariant with respect to S_t , we can write them in terms of $X_t = \ln (S_t/S_0)$ and V_t :

$$dX_{t} = \left(r - \frac{1}{2}V_{t}\right)dt + \sqrt{V_{t}}\left(\rho dB_{t} + \sqrt{1 - \rho^{2}}dW_{t}\right), \quad X_{0} = 0, \quad (7)$$
$$dV_{t} = \Phi\left(V_{t}\right)dt + \Psi\left(V_{t}\right)dB_{t}, \quad V_{0} = v,$$

¹We intend to cover the latter topic in a separate publication.

Alternatively, we can study the joint evolution of the price S_t and its volatility σ_t :

$$\frac{dS_t}{S_t} = rdt + \sigma_t \left(\rho dB_t + \sqrt{1 - \rho^2} dW_t \right), \quad S_0 = s, \\
d\sigma_t = \phi \left(\sigma_t \right) dt + \psi \left(\sigma_t \right) dB_t, \quad \sigma_0 = \sigma.$$
(8)

Given that Eqs (8) are scale invariant with respect to S_t , we can write them in terms of $X_t = \ln (S_t/S_0)$ and σ_t :

$$dX_t = \left(r - \frac{1}{2}\sigma_t^2\right)dt + \sigma_t \left(\rho dB_t + \sqrt{1 - \rho^2} dW_t\right), \quad X_0 = 0, d\sigma_t = \phi\left(\sigma_t\right)dt + \psi\left(\sigma_t\right)dB_t, \quad \sigma_0 = \sigma,$$
(9)

which is often more convenient. From now on, we shall concentrate of studying the dynamics of X_t .

In the general case, we can write dB_t in the form

$$dB_t = \frac{dV_t - \Phi\left(V_t\right)dt}{\Psi\left(V_t\right)},\tag{10}$$

and obtain the following conditionally-independent dynamics for the log-price X_t :

$$dX_t = \left(r - \frac{V_t}{2} - \frac{\rho\sqrt{V_t}\Phi\left(V_t\right)}{\Psi\left(V_t\right)}\right)dt + \frac{\rho\sqrt{V_t}dV_t}{\Psi\left(V_t\right)} + \sqrt{1 - \rho^2}\sqrt{V_t}dW_t, \quad X_0 = 0.$$
(11)

Similarly, we can write dB_t in the form

$$dB_t = \frac{d\sigma_t - \phi(\sigma_t)}{\psi(\sigma_t)},\tag{12}$$

and get the following dynamics for X_t :

$$dX_t = \left(r - \frac{1}{2}\sigma_t^2 - \frac{\rho\sigma\phi\left(\sigma_t\right)}{\psi\left(\sigma_t\right)}\right)dt + \frac{\rho\sigma_t d\sigma_t}{\psi\left(\sigma_t\right)} + \sqrt{1 - \rho^2}\sigma_t dW_t, \quad X_0 = 0.$$
(13)

Assuming that the variance or volatility paths are given, Eqs (11), (13) describe drifted arithmetic Brownian motion with time-dependent drift and volatility.

For the well-known Heston model, Heston (1993), we have

$$\Phi(V_t) = \kappa (\theta - V_t), \quad \Psi(V_t) = \varepsilon \sqrt{V_t}, dV_t = \kappa (\theta - V_t) dt + \varepsilon \sqrt{V_t} dB_t,$$
(14)

so that Eq. (11) has the form

$$dX_t = \left(r - \frac{\rho\kappa\theta}{\varepsilon} - \left(\frac{1}{2} - \frac{\rho\kappa}{\varepsilon}\right)V_t\right)dt + \frac{\rho}{\varepsilon}dV_t + \sqrt{1 - \rho^2}\sqrt{V_t}dW_t.$$
 (15)

Thus, X_t is the so-called drifted Brownian motion driven by the stochastic differential equation of the form

$$dX_t = \mu(t) dt + \nu(t) dW_t,$$

$$\mu(t) = \left(r - \frac{\rho\kappa\theta}{\varepsilon} - \left(\frac{1}{2} - \frac{\rho\kappa}{\varepsilon}\right) V_t\right) + \frac{\rho}{\varepsilon} \frac{dV_t}{dt},$$

$$\nu(t) = \sqrt{1 - \rho^2} \sqrt{V_t}.$$
(16)

For the Stein-Stein model, Stein and Stein (1991), we have

$$\phi(\sigma_t) = \hat{\kappa} \left(\hat{\theta} - \sigma_t \right), \quad \psi(\sigma_t) = \hat{\varepsilon}, \\
d\sigma_t = \hat{\kappa} \left(\hat{\theta} - \sigma_t \right) dt + \hat{\varepsilon} dB_t,$$
(17)

so that Eq. (13) becomes

$$dX_t = \left(r - \frac{\rho \hat{\kappa} \hat{\theta}}{\hat{\varepsilon}} \sigma_t - \left(\frac{1}{2} - \frac{\rho \hat{\kappa}}{\hat{\varepsilon}}\right) \sigma_t^2\right) dt + \frac{\rho}{\hat{\varepsilon}} \sigma_t d\sigma_t + \sqrt{1 - \rho^2} \sigma_t dW_t.$$
(18)

Traditionally, the Heston model is viewed as better describing the market than the Stein-Stein model since the latter does allow for zero variance. In our opinion, this is not a particularly important issue, outweighed by many advantages of the Stein-Stein model, such as a very straightforward way of simulating the evolution of the volatility path. The Heston model gained popularity due to the simple fact that it is exactly solvable for vanilla options. The explicit solution can be calculated via the original Heston (1993) formula. However, the Lewis-Lipton formula is much more efficient; see Lewis (2000), Lipton (2001), Lipton (2002), Lipton and Sepp (2008), and Schmeltze (2010).

In this paper, we consider the Heston model with constant coefficients to follow a long-established tradition, even though it is not necessarily our preferred model. We emphasize that our method is general and can handle any reasonable stochastic volatility model.

2.2 Analytical valuation formula for vanilla options

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The popularity of the Heston model stems from the fact that one can write its solution in the closed-form; see Heston (1993). However, experience has shown that using the original formula is difficult due to several technical drawbacks. Therefore, for benchmarking purposes, here we present the Lewis-Lipton formula, which is easy to implement and use in practice:

$$C^{H}(0, S_{0}, K, T; r, \rho, \kappa, \theta, \varepsilon, V_{0}) = S_{0}\left(1 - \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{e^{(i\chi + 1/2)(\ln(K/S_{0}) - rT) + \alpha(T,\chi) - (\chi^{2} + 1/4)\beta(T,\chi)V_{0}}}{(\chi^{2} + 1/4)} d\chi\right),$$

$$\alpha(T, \chi) = -\frac{\kappa\theta}{\varepsilon^{2}} \left[\psi_{+}T + 2\ln\left(\frac{\psi_{-} + \psi_{+}\exp(-\zeta T)}{2\zeta}\right)\right],$$

$$\beta(T, \chi) = \frac{1 - \exp(-\zeta T)}{\psi_{-} + \psi_{+}\exp(-\zeta T)},$$

$$\psi_{\pm} = \mp (i\rho\varepsilon\chi + \hat{\kappa}) + \zeta,$$

$$\zeta = \sqrt{\varepsilon^{2}} (1 - \rho^{2}) \chi^{2} + 2i\varepsilon\rho\hat{\kappa}\chi + \hat{\kappa}^{2} + \frac{\varepsilon^{2}}{4},$$
(19)

where $\hat{\kappa} = \kappa - \rho \varepsilon / 2$. Further details are given in Lewis (2000), Lipton (2001), Lipton (2002), and Schmeltze (2010).² It is clear that Eq. (19) is a generalization of Eq. (4).

²There is a typo in Lipton (2002) - a minus sign in front of β . This typo is corrected in Lipton (2018), Chapter 10.

2.3 Conditional valuation formula for vanilla options

Unfortunately, with very few exceptions, finding a closed-form solution for barrier or other exotic options on assets with stochastic volatility is not possible, even if such a solution exists for vanilla options. Hence, more general volatility models for barrier options are as good (or bad) as the more traditional Heston and Stein-Stein models, which enjoy closed-form solutions for vanilla options.

We express the log-return process as a linear combination of the two processes:

$$X_t = Y_t + \left(\left(r - \frac{\rho \kappa \theta}{\varepsilon} \right) t - \left(\frac{1}{2} - \frac{\rho \kappa}{\varepsilon} \right) I_t \right) + \frac{\rho}{\varepsilon} \left(V_t - V_0 \right) \equiv Y_t + M_t, \quad (20)$$

where

$$Y_t = \sqrt{1 - \rho^2} \int_0^t \sqrt{V_t} dW_t, \ Y_0 = 0,$$

$$I_t = \int_0^t V_{t'} dt', \ I_0 = 0,$$

$$M_t = \left(\left(r - \frac{\rho \kappa \theta}{\varepsilon} \right) t - \left(\frac{1}{2} - \frac{\rho \kappa}{\varepsilon} \right) I_t \right) + \frac{\rho}{\varepsilon} \left(V_t - V_0 \right), \quad M_0 = 0.$$
(21)

Accordingly, conditionally on the filtration generated by the variance process V_t , we can represent the solution to the price process given by Eq. (6) as follows:

$$S_t = e^{M_t + Y_t} S_0, (22)$$

where M_t is interpreted as a time-deterministic cumulative drift, and Y_t is a martingale with a deterministic time-dependent quadratic variance.

It is clear that pricing for path-independent options, such as European calls and puts, simplifies to the Black-Scholes-Merton formula, provided that either a variance path $\{V_t | 0 \le t \le T\}$ (or a volatility path $\{\sigma_t | 0 \le t \le T\}$) is given. To be concrete, consider a call option with maturity T and strike K. Eq. (6) yields

$$\frac{dS_t}{S_t} = \left(\left(r - \frac{1}{2} V_t \right) dt + \rho \sqrt{V_t} dB_t \right) + \sqrt{1 - \rho^2} \sqrt{V_t} dW_t, \quad , \quad S_0 = s.$$
(23)

Thus,

$$S_T = e^{rT - \frac{1}{2} \left(1 - \rho^2\right) I_T + \sqrt{1 - \rho^2} \sqrt{I_T} \eta} \left(e^{-\frac{1}{2} \rho^2 I_T + \rho J_T} S_0 \right),$$
(24)

where the non-dimensional random variables I_T, J_T , are given by

$$I_T = \int_0^T V_t dt, \quad J_T = \int_0^T \sqrt{V_t} dB_t, \tag{25}$$

and η is the standard (0, 1) normal variable. Accordingly for a particular trajectory $\{V_t | 0 \le t \le T\}$, we obtain the following expression

$$C = C^{BS} \left(0, e^{-\frac{1}{2}\rho^2 I_T + \rho J_T} S_0, T, K; r, \sqrt{1 - \rho^2} \sqrt{\frac{I_T}{T}} \right),$$
(26)

where the values (I_T, J_T) are assumed to be known. The unconditional price is obtained by averaging over all possible (I_T, J_T) :

$$C^{H}(0, S_{0}, K, T; r, \rho, \kappa, \theta, \varepsilon, v_{0}) = \int_{0}^{\infty} \int_{0}^{\infty} C^{BS}\left(0, e^{-\frac{1}{2}\rho^{2}I_{T}+\rho J_{T}}S_{0}, T, K; r, \sqrt{1-\rho^{2}}\sqrt{\frac{I}{T}}\right) \Xi\left(I_{T}, J_{T}\right) dJ_{T} dI_{T},$$
(27)

where $\Xi(I_T, J_T)$ is the joint probability density function (pdf) for the pair (I_T, J_T) ; see Willard (1997) and Romano and Touzi (1997). Thus, Eq. (27) splits the calculation of the call option price into two stages. First, the conditional price is found analytically via the standard Black-Scholes formula. Second, this conditional price is averaged according to the particular choice of the process for the variance V_t . Of course, the first stage is trivial. The usefulness of this formula depends on how easy (or difficult) it is to find the pdf for (I, J) and complete the second stage.

Two approaches have been used in practice - the standard Monte-Carlo method for calculating $\{V_t | 0 \le t \le T\}$, I, J, and a more advanced (but much harder) method based on the augmentation principle described in Section (13.2) Lipton (2001). Specifically, the augmented dynamic equation for V_t yields the following system of degenerate PDEs for the triple (V, I, J):

$$dV_{t} = \Phi(V_{t}) dt + \Psi(V_{t}) dB_{t}, \quad V_{0} = v, dI_{t} = V_{t} dt, \quad I_{0} = 0, dJ_{t} = \sqrt{V_{t}} dB_{t}, \quad J_{0} = 0.$$
(28)

The corresponding Green's function G(t, V, I, J) is governed by the degenerate Fokker-Planck equation of the form

$$G_{t} + (\Phi(V)G)_{V} + VG_{I} - \frac{1}{2} (\Psi^{2}(V)G)_{VV} - (\sqrt{V}\Psi(V)G)_{VJ} - \frac{1}{2} (VG)_{JJ} = 0,$$

$$G(0, V, I, J) = \delta(V - v) \delta(I) \delta(J).$$
(29)

In general, solving this equation is complicated; however, for the Heston model, Lipton (2001) found a closed-form solution.

In Figure 1 we compare prices of European call options given by analytical formula (19), and Monte Carlo formula (27). As expected, both formulas agree, although the latter formula is much slower than the former.

3 Barrier options

3.1 Formulation

Our task is to price a barrier option written on an asset with stochastic volatility. For brevity, we consider barrier options with the lower barrier $B < S_0$ only. Considering other possibilities, such as pricing options with the upper barrier or popular double-no-touch options, is left to the reader. The corresponding



Figure 1: Implied volatilities of European call options. We obtain the corresponding prices by using Eqs. (19) and (27). Here, and throughout the paper we use the following parameters: $S_0 = 1$, $V_0 = 0.25$, T = 1.0, r = 0.03, $\kappa = 1.0$, $\theta = 0.2$, $\rho = -0.3$, $\varepsilon = 0.4$.

IBVP can be written in the form

$$P_{t} + (r - \frac{1}{2}V) P_{X} + \kappa (\theta - V) P_{V} + \frac{1}{2}V (P_{XX} + 2\rho P_{XV} + P_{VV}) - rP = 0,$$

$$P (T, X, V) = \Pi (X),$$

$$P (t, \xi, V) = 0,$$
(30)

where $\xi = \ln (B/S_0) < 0$, $\Pi(X)$ is the terminal payoff function. Typical examples include the no-touch, call and put payoffs defined by:

$$\Pi(X_T) = 1, \ \Pi(X_T) = S_0 \left(e^{X_T} - e^k \right)_+, \ \Pi(X_T) = S_0 \left(e^k - e^{X_T} \right)_+, \quad (31)$$

where $k = \ln (K/S_0)$. Alternatively, we can write the adjoint problem for Green's function G:

$$G_{t} + \left(\left(r - \frac{1}{2}V \right) G \right)_{X} + \left(\kappa \left(\theta - V \right) G \right)_{V} - \frac{1}{2} \left(\left(VG \right)_{XX} + 2\rho \left(VG \right)_{XV} + \left(VG \right)_{VV} \right) + rG = 0, G \left(0, X, V \right) = \delta \left(X \right) \delta \left(V - V_{0} \right), G \left(t, \xi, V \right) = 0,$$
(32)

Once the corresponding Green's function is calculated, we can find P via simple integration:

$$P(0,0,V_0) = \int_{\xi}^{\infty} \int_0^{\infty} G(0,0,V_0,T,X_T,V_T) \Pi(X_T) \, dV_T dX_T.$$
(33)

While such options can be priced via either FDM or MCS, both are notoriously slow. Therefore, we want to design a much faster method, enjoying equal or higher accuracy than the classical alternatives.

As far as analytical solutions are concerned, only one is known. It was discovered by Lipton (1997), see also Lipton (2001), Lipton and McGhee (2002), and Andreasen (2001). Lipton (1997) observed that in the special case $r = 0, \rho = 0$, IBVPs (30), (32) are symmetric with respect to the transformation $X \to -X$. Hence, the classical method of images is applicable, and solutions to these problems can be presented as a linear combination of solutions without barriers. Of course, one can use this approach for options in the presence of an upper barrier, as well as for double-barrier options.

Recently, there were several unsuccessful attempts to solve the pricing problem with $r^2 + \rho^2 > 0$. For example, De Gennaro Aquino and Bernard (2019) presented a solution, relying on an explicit expression for the joint distribution of the value of a Brownian motion with time-dependent drift and its maximum and minimum; it was quickly shown by one of the present authors that their calculation is erroneous; see De Gennaro Aquino and Bernard (2021).³ He and Lin (2021) presented a "solution", which relies on the unsubstantiated replacement of the time-dependent drift by a constant. Their approach is so arbitrary and frivolous that its detailed repudiation is not warranted.

 $^{^{3}}$ Finding the joint distribution for a Brownian motion with time-dependent drift and its maximum and minimum is challenging. We present its solution in Section 5.

3.2 Conditional valuation formula for barrier options

It is hard to extend interesting formula (27) for barrier options. However, it is not impossible! Following Section 13.3 in Lipton (2001), we express the value of a path-dependent option as an integral in the functional space of price trajectories:

$$P(S_0) = e^{-rT} \int_{\Omega} \mathcal{F}(\omega) d\mathcal{D}(\omega)$$
(34)

where $\mathcal{F}(\omega)$ is a functional mapping of the space of trajectories into payoffs, and $\mathcal{D}(\omega)$ is the risk-neutral probability measure.

Further, by applying the augmentation principle, we introduce the functional Λ_t to represent the path-dependent variable linked to the evolution of the spot price S_t . We then consider evaluation of a derivatives security with the terminal pay-off function $f(S, \Lambda)$. Finally, we extend the joint dynamics in Eq. (6) with the dynamics of augmented variable $\Lambda_t = \min_{0 \le t' \le t} S_t$:

$$S_t = rS_t dt + \sqrt{V_t} S_t \left(\rho dB_t + \sqrt{1 - \rho^2} dW_t \right), \quad S_0 = s,$$

$$dV_t = \Phi(V_t) dt + \Psi(V_t) dB_t, \quad V_0 = v,$$

$$d\Lambda_t = \theta \left(\Lambda_t - S_t \right) \left(dS_t \right)_-, \quad \Lambda_0 = S_0.$$
(35)

The payoff function $f(S, \Lambda)$ is given by:

$$f(S_T, \Lambda_T) = \mathbf{1}_{\{\Lambda_T > B\}} \Pi(S_T)$$
(36)

The joint dynamic, given by Eqs (35), is Markovian. Accordingly, we can reduce the general formula (34) to the form:

$$P(0,S_0) = e^{-rT} \int_{V=0}^{\infty} \int_{S=0}^{\infty} \int_{\Lambda=0}^{\infty} f(S,\Lambda) G(T,V,S,\Lambda;0,V_0,S_0,\Lambda_0) d\Lambda dS dV.$$
(37)

Here $G(T, V, S, \Lambda; 0, V_0, S_0, \Lambda_0)$ is the risk-neutral probability density function for the joint evolution of state variables in SDE (35). We emphasize that while the payoff function does not depend on the variance V, the valuation problem has three spatial variables, including variance, in addition to the time variable.

Below, we introduce a novel method to solve the valuation problem (37) semi-analytically. Using Eq. (22), the density of the price S_t is log-normal with time-dependent drift and variance conditional on the filtration generated by stochastic variance V_t , \mathbb{F}^V . Therefore, we represent the valuation formula (37) as follows:

$$P(0,S_0) = e^{-rT} \mathbb{E}_V \left[\int_{S=0}^{\infty} \int_{\Lambda=0}^{\infty} f(S,\Lambda) \Gamma\left(T,S,\Lambda;0,S_0,\Lambda_0 | V(\omega)\right) d\Lambda dS \right],$$
(38)

where Γ is the Gaussian density of S_t and Λ_t conditional of the variance path $V(\omega)$, and the expectation is computed over all paths of variance process V_t . The term in the square brackets is the value of a path-dependent option on Λ and S, computed in closed-form or numerically.

We apply the MHP to compute the inner integral for barrier options in the closed-form. Thus, we have generalized Eq (27), valid for path-independent options, to path-dependent options, and apply the new result to value barrier options in the semi-closed form.

Mixed MHP-MCS approach to solving IBVPs 4

We consider the IBVP (30). In the spirit of Section 3.2, we split our algorithm in two steps: the outer MCS loop, which generates a bunch of trajectories for the stochastic variance v_t , and the inner loop, which solves the one-dimensional IBVP for X_t . We discuss two approaches for solving the latter problem: the finite difference approach developed by Loeper and Pironneau (2009), and the MHP approach inspired by Lipton *et al.* (2019), Lipton and Kaushansky (2020). We start with the inner loop. The outer loop is relatively straightforward. It requires to average the inner price by using the equation for V_t . This will provide solution of the form given by Eq. (2) in Lipton and McGhee (2002). The simplest way of doing this averaging is via MCS, although in some exceptional cases other approaches can be envisaged.

4.1**One-dimensional IBVP**

Eq. (15), conditional on the variance path, can be written as an SDE of the form

$$dX_t = \mu(t) dt + \nu(t) dW_t, \quad X_t \ge \xi, \tag{39}$$

where X_t the drifted Brownian motion with μ, ν given by Eqs (16).

While studying X_t on the entire axis $(-\infty, \infty)$ is almost trivial, dealing with the same process in a semi-bounded domain (ξ, ∞) might be quite hard. This paper shows how to do it by using the FDM and the MHP.

We can always rescale time t by introducing v, such that dv = v(t) dt. As a result, the governing SDE becomes

$$dX_{\upsilon} = \lambda\left(\upsilon\right)d\upsilon + dW_{\upsilon},\tag{40}$$

where $\lambda(v) = \mu(t) / \nu(t), v(t) = \int_0^t v(t') dt'$. Using the decomposition for X_t given by Eq. (20) we present the valuation problems follows:

$$dX_t = dY_t + dM_t, \quad Y_0 = 0, \quad M_0 = 0, \quad X_t \ge \xi,$$
(41)

where Y_t is stochastic part and M_t is deterministic part, and $\xi < 0$. Then we can express the problem (20) in terms of stochastic variable Y_t as follows:

$$dY_t = \sqrt{1 - \rho^2} \sqrt{V_t} dW_t, \quad Y_0 = 0, \quad Y_t \ge \xi - M_t.$$
 (42)

We use a new time variable

$$v(t) = (1 - \rho^2) I_t,$$
 (43)

and write

$$dY_{\upsilon} = dW_{\upsilon}, \quad Y_0 = 0, \quad Y_{\upsilon} \ge \xi - N_{\upsilon},$$
(44)

where

$$N_{\upsilon} = M_{t(\upsilon)}.\tag{45}$$

Thus, we have one of the two venues to explore: (A) studying the processes X_t or X_v given by Eq. (39) and Eq. (40), respectively; (B) dealing with the processes Y_t or Y_v given by Eq. (42) and Eq. (44). In case (A) there is non-zero time-dependent drift and flat boundary; in case (B) there is no drift but the boundary is time-dependent.

Given a variance trajectory, we need to discuss how to calculate μ, ν, λ, M , and N. Let $\{v_k | k = 0, 1, ..., K\}$, $v_0 = v$, be a particular path generated via discretization of Eqs (14) with homogeneous time-step $\Delta t = T/K$. This equation can be discretized in various ways, for example, via the Euler-Maryama scheme or the Milstein scheme. For special cases such as the Feller process corresponding to the Heston model, there are clever schemes tailored to the specific process at hand. However, we are not pursuing them here since it is unnecessary to achieve our objective. We have

$$v_{0} = v, \quad v_{k} = v_{k-1} + \kappa \left(\theta - v_{k-1}\right) \Delta t + \varepsilon \sqrt{\Delta t} \eta, \\ I_{0} = 0, \quad I_{k} = I_{k-1} + \frac{\Delta t}{2} \left(v_{k} + v_{k-1}\right),$$
(46)

$$\mu_{k} = r - \frac{\rho \kappa \theta}{\varepsilon} - \left(\frac{1}{2} - \frac{\rho \kappa}{\varepsilon}\right) v_{k} + \frac{\rho}{\varepsilon} \frac{\left(\left(v_{k} - v_{k-1}\right)\right)}{\Delta t},$$

$$\nu_{k} = \sqrt{1 - \rho^{2}} \sqrt{v_{k}},$$

$$M_{0} = 0, \quad M_{k} = \frac{kT}{K} \left(r - \frac{\rho \kappa \theta}{\varepsilon}\right) - \left(\frac{1}{2} - \frac{\rho \kappa}{\varepsilon}\right) I_{k} + \frac{\rho}{\varepsilon} \left(v_{k} - v\right).$$

$$(47)$$

It is clear that

$$\upsilon_k = \left(1 - \rho^2\right) I_k, \quad \Upsilon = \left(1 - \rho^2\right) I_K, \tag{48}$$

where $\{v_k | k = 0, 1, ..., K\}$ is an inhomogeneous partition of the interval $[0, \Upsilon]$. For our purposes, we treat the sequences $\{\lambda_k = \mu_k / \nu_k | k = 0, 1, ..., K\}$ and $\{M_k | k = 0, 1, ..., K\}$ as functions of v_k , which is possible because v_k is a monotonically increasing sequence, and interpret them accordingly.

We illustrate the corresponding functions in Figures 2, 3. We emphasize that μ and λ are very irregular, since they depend on the white noise process dW_t/dt , so that we have to deal with random terms of order $\Delta t^{-1/2}$. At the same time, the moving boundary M_t is much more regular, and depends on random terms of order $\Delta t^{1/2}$.

4.2 Solving IBVPs via Crank-Nicolson method

Let us describe how to price a conditional barrier option via the FDM; see Loeper and Pironneau (2009).⁴ Specifically, we want to solve the following

 $^{^4 {\}rm Surprisingly},$ Loeper and Pironneau ~(2009) do not comment on the irregular nature of the drift coefficient.



Figure 2: Time-dependent coefficients (a) advection, (b) diffusion. The corresponding parameters are the same as in Figure 1. Here $N_t = 52$ - one step per week.







Figure 3: (a) $t(\Upsilon)$, (b) $\lambda(\Upsilon)$

.

problem on the semi-axis (ξ, ∞) :

$$P_{t}(t,X) + \mu(t) P_{X}(t,X) + \frac{1}{2}\nu(t) P_{XX}(t,X) - \varkappa(t) P(t,X) = 0,$$

$$P(T,X) = \Pi(X), \quad P(t,\xi) = 0.$$
(49)

We introduce \bar{P} , such that

$$P(t,X) = e^{-\int_t^T \varkappa(t') dt'} \bar{P}(t,X), \qquad (50)$$

and get the following problem for $\bar{P}(t, X)$:

$$\bar{P}_{t}(t,X) + \mu(t)\bar{P}_{X}(t,X) + \frac{1}{2}\nu(t)\bar{P}_{XX}(t,X) = 0,
\bar{P}(T,X) = \Pi(X), \quad \bar{P}(t,\xi) = 0.$$
(51)

In the context we are interested in, $\varkappa(t) = r$, so that $P = \exp(r(t-T))\overline{P}$.

We choose a uniform spatial grid $\{x_m = \xi + m (U - \xi) / M | m = 0, 1, ..., M\}$, where U is a sufficiently remote upper boundary such that 0 belongs to the spatial grid, $x_o = 0$, and a temporal grid $\{t^n | t^0 < t^1 < ... < t^N = T\}$, which is not necessarily uniform. Then, we apply the usual Crank-Nicolson method for the advection-diffusion diffusion and get the following system of matrix equations for \bar{P}_m^n :

$$\bar{P}_{m}^{n+1} - \bar{P}_{m}^{n} + \alpha^{n+1/2} \left(\bar{P}_{m+1}^{n+1} - \bar{P}_{m-1}^{n+1} + \bar{P}_{m+1}^{n} - \bar{P}_{m-1}^{n} \right)
+ \beta^{n+1/2} \left(\bar{P}_{m+1}^{n+1} - 2\bar{P}_{m}^{n+1} + \bar{P}_{m-1}^{n+1} + \bar{P}_{m+1}^{n} - 2\bar{P}_{m}^{n} + \bar{P}_{m-1}^{n} \right) = 0.$$
(52)

where

$$\alpha^{n+1/2} = \frac{\mu^{n+1/2} \Delta t^{n+1/2}}{4\Delta x}, \quad \beta^{n+1/2} = \frac{\nu^{n+1/2} \Delta t^{n+1/2}}{4\Delta x^2}.$$
 (53)

We emphasize that given the extreme irregularity of μ , using more complicated approaches for treading the drift term is not warranted.

We can write the system of equations in matrix form.

$$\mathbb{A}^{n+1/2}_{(\alpha,\beta)}\bar{P}^n = \mathbb{A}^{n+1/2}_{(-\alpha,-\beta)}\bar{P}^{n+1}, \\ \bar{P}^N_m = \Pi(x_m),$$
(54)

where

For brevity, the superscripts are omitted. At infinity we apply one-sided derivatives and use the PDE *itself* to formulate the boundary condition.

Alternatively, we can rescale time, reduce the problem (29) to the form:

$$\bar{P}_{v} + \lambda (v) \bar{P}_{x} + \frac{1}{2} \bar{P}_{xx} = 0, \bar{P} (\Upsilon, x) = \Pi (x), \quad \bar{P} (v, \xi) = 0,$$
(56)

and apply the Crank-Nicolson method to problem (56).

Of course, we can attack the pricing problem by solving the corresponding Fokker-Planck equation for Green's function G:

$$G_{t}(t,X) + \mu(t) G_{X}(t,X) - \frac{1}{2}\nu(t) G_{XX}(t,X) + \varkappa(t) G(t,X) = 0,$$

$$G(0,X) = \delta(X), \quad G(t,\xi) = 0,$$
(57)

and write

$$P(0,0) = \int_{0}^{\infty} G(0,0;T,x) \Pi(x) \, dx.$$
(58)

We introduce \overline{G} , such that

$$G(t,X) = e^{-\int_0^t \varkappa(t') dt'} \bar{G}(t,X).$$
(59)

Accordingly,

$$\bar{G}_{t}(t,X) + \mu(t)\bar{G}_{X}(t,X) - \frac{1}{2}\nu(t)\bar{G}_{XX}(t,X) = 0, \bar{G}(0,X) = \delta(X), \quad \bar{G}(t,\xi) = 0,$$
(60)

$$P(0,0) = e^{-\int_0^T \varkappa(t')dt'} \int_0^\infty \bar{G}(0,0;T,x) \Pi(x) \, dx.$$
(61)

As before, we apply the Crank-Nicolson method for the advection-diffusion and get the following system of matrix equations for \bar{G}_m^n :

$$\bar{G}_{m}^{n+1} - \bar{G}_{m}^{n} + \alpha^{n+1/2} \left(\bar{G}_{m+1}^{n+1} - \bar{G}_{m-1}^{n+1} + \bar{G}_{m+1}^{n} - \bar{G}_{m-1}^{n} \right)
-\beta^{n+1/2} \left(\bar{G}_{m+1}^{n+1} - 2\bar{G}_{m}^{n+1} + \bar{G}_{m-1}^{n+1} + \bar{G}_{m+1}^{n} - 2\bar{G}_{m}^{n} + \bar{G}_{m-1}^{n} \right) = 0.$$
(62)

In the matrix form we get

$$\mathbb{A}^{n+1/2}_{(\alpha,\beta)}\bar{G}^{n+1} = \mathbb{A}^{n+1/2}_{(-\alpha,-\beta)}\bar{G}^n, \qquad (63)$$
$$\bar{G}^0 = \delta_{o,m}.$$

Once \bar{G}^N is found, we can represent P_k^0 as

$$P_k^0 = e^{-rT} \Delta x \sum_{m=1}^M \bar{G}_m^N \Pi_m.$$
 (64)

4.3 Solving IBVPs via MHP

4.3.1 Analytic results

We wish to find Green's function for process (44), or, equivalently, to solve the following IBVP in a bounded domain with a moving boundary:

$$\frac{\partial}{\partial v}\bar{G}(v,Y) = \frac{1}{2}\frac{\partial^2}{\partial Y^2}\bar{G}(v,Y),
\bar{G}(0,Y) = \delta(Y), \quad \bar{G}(v,\xi - N_v) = 0, \quad \bar{G}(v,Y \to \infty) \to 0,
\xi - N_v \le Y < \infty, \quad 0 \le v \le \Upsilon.$$
(65)



Figure 4: A typical lower boundary $\xi - N_v$ as a function of v.

For a representative Monte Carlo path, the corresponding boundary is shown in Figure 4.

We split \overline{G} , and represent it in the form:

$$\bar{G}(v,Y) = H(v,Y) - F(v,Y), \qquad (66)$$

where H(v, Y) is the standard heat kernel,

$$H(v,Y) = \frac{\exp\left(-\frac{Y^2}{2v}\right)}{\sqrt{2\pi v}},\tag{67}$$

and F(v, Y) solves the following problem:

$$\frac{\partial}{\partial v}F(v,Y) = \frac{1}{2}\frac{\partial^2}{\partial Y^2}F(v,Y), \quad \xi - N_v \le Y < \infty, F(0,Y) = 0, \quad F(v,\xi - N_v) = f(v), \quad F(v,Y \to \infty) \to 0,$$
(68)

where

$$f(\upsilon) = H\left(\upsilon, \xi - N_{\upsilon}\right). \tag{69}$$

The MHP allows one to represent F(v, Y) in the form

$$F(v,Y) = \int_0^v \frac{(Y - \xi + N_{v'}) \exp\left(-\frac{(Y - \xi + N_{v'})^2}{2(v - v')}\right)}{\sqrt{2\pi (v - v')^3}} \phi(v') \, dv', \qquad (70)$$

where $\phi(v)$ solves the Volterra equation of the second kind:

$$\phi(v) + \int_0^v \frac{\Theta(v, v') \Xi(v, v')}{\sqrt{2\pi (v - v')}} \phi(v') \, dv' = f(v) \,, \tag{71}$$

and

$$\Theta(v,v') = -\frac{N_v - N_{v'}}{(v - v')}, \quad \Xi(v,v') = \exp\left(-\frac{(v - v')\Theta^2(v,v')}{2}\right), \quad (72)$$
$$\Theta(v,v) = -\frac{dN_v}{dv}, \quad \Xi^>(v,v) = 1.$$

Assuming that $\phi(v)$ is known, we can represent $\overline{G}(\Upsilon, Y)$ as follows:

$$\bar{G}(\Upsilon, Y) = H(\Upsilon, Y) - F(\Upsilon, Y), \qquad (73)$$

where $F(\Upsilon, Y)$ is given by Eq. (70). Finally, returning back to the original variables, we get

$$\bar{G}(T,X) = H\left(\Upsilon\left(T\right), X - M_T\right) - F\left(\Upsilon\left(T\right), X - M_T\right).$$
(74)

When $Z = X - \xi \to 0$, the corresponding integral has to be dealt with carefully due to a singularity at v' = v. We have

$$F(v, X - N_v) = \int_0^v \frac{(Z - N_v + N_{v'}) \exp\left(-\frac{(Z - N_v + N_{v'})^2}{2(v - v')}\right)}{\sqrt{2\pi(v - v')^3}} \phi(v') dv'$$

= $e^{ZN'(v)} \phi(v) \int_0^v \frac{Z \exp\left(-\frac{Z^2}{2(v - v')}\right)}{\sqrt{2\pi(v - v')^3}} dv' + \int_0^v \frac{I^{(1)}(v,v')}{\sqrt{(v - v')}} dv' + \int_0^v \frac{I^{(2)}(v,v')}{\sqrt{(v - v')}} dv'$
= $2e^{ZN'(v)} \phi(v) \Re\left(-\frac{Z}{\sqrt{v}}\right) + \int_0^v \frac{I^{(1)}(v,v')}{\sqrt{(v - v')}} dv' + \int_0^v \frac{I^{(2)}(v,v')}{\sqrt{(v - v')}} dv',$ (75)

where

$$I^{(1)}(v,v') = \frac{Z \exp\left(-\frac{Z^2}{2(v-v')}\right) \left(\exp(-Z\Theta(v,v'))\Xi(v,v')\phi(v') - e^{ZN'(v)}\phi(v)\right)}{\sqrt{2\pi}(v-v')}, \quad (76)$$
$$I^{(2)}(v,v') = \frac{\Theta(v,v') \exp\left(-\frac{(Z-N_v+N_{v'})^2}{2(v-v')}\right)\phi(v')}{\sqrt{2\pi}}.$$

It is clear that integrals in Eq. (75) have weak singularities and hence are easy to handle.

4.3.2 Numerics

There are numerous well-known approaches to solving Volterra equations; see, Linz (1985), among many others. We choose the most straightforward approach and show how to solve the following archetypal Volterra equation with weak singularity numerically:

$$\phi(\upsilon) + \int_0^{\upsilon} \frac{K(\upsilon, \upsilon')}{\sqrt{\upsilon - \upsilon'}} \phi(\upsilon') \, d\upsilon' = f(\upsilon), \tag{77}$$

where K(v, v') is a non-singular kernel. We write

$$\int_{0}^{\upsilon} \frac{K(\upsilon, \upsilon')\phi(\upsilon')}{\sqrt{\upsilon - \upsilon'}} d\upsilon' = -2 \int_{0}^{\upsilon} K(\upsilon, \upsilon')\phi(\upsilon') \ d\sqrt{\upsilon - \upsilon'}.$$
 (78)

We map this equation to a grid $0 = v_0 < v_1 < \ldots < v_N = v$. We introduce the following notation:

$$f_{k} = f(v_{k}), \quad \phi_{k} = \phi(v_{k}), \quad K_{k,l} = K(v_{k}, v_{l}), \\ \Delta_{k,l} = v_{k} - v_{l}, \quad \Pi_{k,l} = \frac{\Delta_{l,l-1}}{(\sqrt{\Delta_{k,l-1}} + \sqrt{\Delta_{k,l}})}.$$
(79)

Then, Eq. (77) can be approximated by the trapezoidal rule as

$$\phi_k + \sum_{l=1}^k \Pi_{k,l} \left(K_{k,l} \phi_l + K_{k,l-1} \phi_{l-1} \right) = f_k, \tag{80}$$

so that

$$\phi_k = \frac{\left(f_k - \sqrt{\Delta_{k,k-1}}K_{k,k-1}\phi_{k-1} - \sum_{l=1}^{k-1}\Pi_{k,l}\left(K_{k,l}\phi_l + K_{k,l-1}\phi_{l-1}\right)\right)}{\left(1 + \sqrt{\Delta_{k,k-1}}K_{k,k}\right)}.$$
 (81)

Thus, ϕ_k can be found by induction starting with $\phi_0 = f_0$. After ϕ_k are determined, F_k can be written in the form

$$F_{k} = 2e^{ZN'_{k}}\phi_{k}\mathfrak{N}\left(-\frac{Z}{\sqrt{\upsilon_{k}}}\right) + \sum_{l=1}^{k}\Pi_{k,l}\left(I^{(1)}_{k,l} + I^{(1)}_{k,l-1}\right) + \sum_{l=1}^{k}\Pi_{k,l}\left(I^{(2)}_{k,l} + I^{(2)}_{k,l-1}\right),$$
(82)

4.3.3Example

Let us consider the special case of constant drift λ ; the corresponding boundary is linear, $\xi - \lambda v$, where v = t. Then Eq. (71) becomes

$$\phi(\upsilon) - \lambda \int_0^{\upsilon} \frac{\exp\left(-\frac{\lambda^2(\upsilon-\upsilon')}{2}\right)}{\sqrt{2\pi\left(\upsilon-\upsilon'\right)}} \phi(\upsilon') \, d\upsilon' = \frac{e^{-\frac{(\xi-\lambda\upsilon)^2}{2\upsilon}}}{\sqrt{2\pi\upsilon}}.$$
(83)

Lipton and Kaushansky (2020) show that

$$\phi(\upsilon) = \frac{e^{-\frac{(\xi - \lambda\upsilon)^2}{2\upsilon}}}{\sqrt{2\pi\upsilon}} + \lambda e^{2\xi\lambda} \Re\left(\frac{\xi + \lambda\upsilon}{\sqrt{\upsilon}}\right),\tag{84}$$

$$F(\upsilon, Y) = \frac{\exp\left(2\xi\lambda - \frac{(Y-2\xi)^2}{2\upsilon}\right)}{\sqrt{2\pi\upsilon}}.$$
(85)

Accordingly,

$$\bar{G}(T,X) = H(T,X-\lambda T) - e^{2\xi\lambda}H(T,X-\lambda T-2\xi).$$
(86)

It is easy to see that $\overline{G}(T,\xi) = 0$, as it should. Figure 5 illustrates our findings.



Figure 5: $\phi(v)$ computed analytically and by solving the Volterra equation. The difference between the corresponding functions is less that $10^{(-4)}$. The parameters are $\xi = -0.5$, $\lambda = 0.5$.

4.3.4 Semi-analytical solution of pricing problems

Green's function is given by Eq. (74). In Figure 6 we compare Green's functions obtained via the FDM and the MHP. The figure shows that these functions are reassuringly close. The MHP is clearly preferred since it is much faster.

Once $\overline{G}(T, X)$ is found all barrier problems can be solved analytically.

For instance, we can calculate the $P(0,T;\xi)$ of the no-touch option for the barrier level $\xi < 0$:

$$P(0,T;\xi) = e^{-rT} \int_{\xi}^{\infty} \left(H\left(\Upsilon, X' - N_{\Upsilon}\right) - F\left(\Upsilon, X' - N_{\Upsilon}\right) \right) dX'$$

$$= e^{-rT} \left(\Re\left(\frac{-(\xi - N_{\Upsilon})}{\sqrt{\Upsilon}}\right) - \int_{0}^{\Upsilon} \frac{\exp\left(-\frac{(\xi - N_{\Upsilon} + N_{U'})^{2}}{2(\Upsilon - \upsilon')}\right)}{\sqrt{2\pi(\Upsilon - \upsilon')}} \phi\left(\upsilon'\right) d\upsilon' \right).$$
(87)

The relative price $C(0, T, K; \xi) / e^{-rT} S_0$ of the barrier call struck at $K = S_0 e^k$, where $k \ge \xi$, has the form:



Figure 6: G(T, x) computed via the FDM and the MHP. The absolute difference between the corresponding functions is less that 10^{-3} . The parameters are the same as in Figure 1. The log-barrier is $\xi = -0.5$.

$$\frac{C(0,S_0,T,K;\xi)}{e^{-r^T S_0}} = \int_k^\infty \left(H\left(\Upsilon, X' - N_\Upsilon\right) - F\left(\Upsilon, X' - N_\Upsilon\right) \right) \left(e^{X'} - e^k \right) dX' \\
= e^{N_\Upsilon} \int_{k-N_\Upsilon}^\infty \frac{e^{-\frac{\eta^2}{2\Upsilon} + \eta}}{\sqrt{2\pi\Upsilon}} d\eta - e^k \int_{k-N_\Upsilon}^\infty \frac{e^{-\frac{\eta^2}{2\Upsilon}}}{\sqrt{2\pi\Upsilon}} d\eta \\
-e^{N_\Upsilon} \int_0^\Upsilon \int_{k-N_\Upsilon}^\infty \frac{(\eta - \xi + N_{v'}) \exp\left(- \frac{(\eta - \xi + N_{v'})^2}{2(\Upsilon - v')^3} + \eta \right)}{\sqrt{2\pi(\Upsilon - v')^3}} d\eta \phi (v') dv' \\
+e^k \int_0^\Upsilon \int_{k-N_\Upsilon}^\infty \frac{(\eta - \xi + N_{v'}) \exp\left(- \frac{(\eta - \xi + N_{v'})^2}{2(\Upsilon - v')^3} \right)}{\sqrt{2\pi(\Upsilon - v')^3}} d\eta \phi (v') dv' \\
= e^{N_\Upsilon + \frac{\Upsilon}{2}} \Re \left(\frac{N_\Upsilon + \Upsilon}{\sqrt{\Upsilon}} \right) - e^k \Re \left(\frac{N_\Upsilon}{\sqrt{\Upsilon}} \right) \\
-e^k \int_0^\Upsilon \frac{\exp\left(- \frac{(k - N_\Upsilon - \xi + N_{v'})^2}{2(\Upsilon - v')} \right)}{\sqrt{2\pi(\Upsilon - v')}} \phi (v') dv' \\
+e^k \int_0^\Upsilon \frac{\exp\left(- \frac{(k - N_\Upsilon - \xi + N_{v'})^2}{2} \right)}{\sqrt{2\pi(\Upsilon - v')}} \phi (v') dv'.$$
(88)

We found that it more efficient to price these options using the backward induction. For example, to price the no-touch option backward, we introduce the new time variable $\varpi = \Upsilon - v$, and the boundary $O_{\varpi} = N_{\Upsilon-v}$, and write

 $P\left(0,T;\xi\right)$ in the form

$$P(0,T;\xi) = e^{-rT} (1 - Q(\Upsilon;\xi)).$$
(89)

Here

$$Q\left(\Upsilon;\xi\right) = \int_{0}^{\Upsilon} \frac{\left(-\xi + O_{\varpi'}\right) \exp\left(-\frac{\left(-\xi + O_{\varpi'}\right)^{2}}{2(\Upsilon-\varpi')}\right)}{\sqrt{2\pi \left(\Upsilon-\varpi'\right)^{3}}} \psi^{(NT)}\left(\varpi'\right) d\varpi', \qquad (90)$$

where $\psi^{(NT)}(\varpi)$ solves Eq. (71) with $f(\varpi) = 1$.

By the same token, we can represent $C(0, S_0, T, K; \xi)$ as follows

$$C(0, S_0, T, K; \xi) = e^{-rT} S_0 \left(D\left(\Upsilon, k\right) - E\left(\Upsilon, k; \xi\right) \right), \tag{91}$$

where

$$D(\Upsilon, k) = e^{k} \left(e^{O_{0}-k+\frac{\Upsilon}{2}} \mathfrak{N} \left(\frac{O_{0}-k+\Upsilon}{\sqrt{\Upsilon}} \right) - \mathfrak{N} \left(\frac{O_{0}-k}{\sqrt{\Upsilon}} \right) \right),$$

$$E(\Upsilon, k; \xi) = \int_{0}^{\Upsilon} \frac{(-\xi+O_{\varpi'}) \exp\left(-\frac{(-\xi+O_{\varpi'})^{2}}{2(\Upsilon-\varpi')^{3}} \right)}{\sqrt{2\pi(\Upsilon-\varpi')^{3}}} \psi^{(C)}(\varpi') d\varpi',$$
(92)

and $\psi^{(C)}(\varpi')$ solves Eq. (71) with

$$f(\varpi) = e^k \left(e^{\xi - O_{\varpi'} + O_0 - k + \frac{\Upsilon}{2}} \mathfrak{N} \left(\frac{\xi - O_{\varpi'} + O_0 - k + \Upsilon}{\sqrt{\Upsilon}} \right) - \mathfrak{N} \left(\frac{\xi - O_{\varpi'} + O_0 - k}{\sqrt{\Upsilon}} \right) \right).$$
(93)

In Figures 7, 8 we compare prices of no-touch and call options obtained via the FDM and the MHP. The figure shows that the corresponding prices are very close. As before, the MHP is clearly much faster than the FDM.

4.4 External loop: averaging over all variance paths

After the corresponding solution is found and expressed in the original variables, we produce a set of random sequences $\{v_k | k = 0, 1, ..., K\}$ and repeat steps. Once a sufficiently large number of paths is generated, we perform averaging and obtain the solution.

In Figures 9, 10, 11, we show the averaged value of Green's function, as well as prices of the no-touch and call options.

5 Joint probability density for the value of drifted Brownian motion and its running minimum

This section, which serves as a mathematical aside, shows that the MHP allows solving a very complex problem of finding the joint probability distribution for a Brownian motion with time-dependent drift and volatility and its



Figure 7: P(T, x) computed via the FDM and the MHP. The absolute difference between the corresponding functions is less that 10^{-3} . The parameters are the same as in Figure 1. The log-barrier is $\xi = -0.5$.



Figure 8: C(T, x) computed via the FDM and the MHP. The absolute difference between the corresponding functions is less that 10^{-4} . The parameters are the same as in Figure 1. The log-barrier is $\xi = -0.5$.



(b)

Figure 9: Green's function averaged over Monte Carlo paths. The number of MC path for the MHP and FDM is 10,000; the number of path for the classical MCS is 100,000.



Figure 10: Price of the no-touch with the log-barrier $= \xi = -0.5$: figure (a), table (b). The number of MC paths for the MHP and MFD is 10,000; the number of patha for the MCS is 100,000.



Strike	MC	Volterra	Pde	Vanilla
0.9000	0.1121	0.1117	0.1115	0.2466
0.9500	0.1040	0.1036	0.1035	0.2215
1.0000	0.0961	0.0957	0.0956	0.1985
1.0500	0.0884	0.0880	0.0880	0.1776
1.1000	0.0811	0.0807	0.0806	0.1586
1.1500	0.0740	0.0736	0.0736	0.1415
1.2000	0.0674	0.0670	0.0670	0.1260
1.2500	0.0612	0.0608	0.0608	0.1122
1.3000	0.0554	0.0551	0.0551	0.0998
1.3500	0.0501	0.0498	0.0497	0.0887
1.4000	0.0452	0.0449	0.0449	0.0789
1.4500	0.0407	0.0404	0.0404	0.0701
1.5000	0.0366	0.0364	0.0364	0.0623

(b)

Figure 11: (a) The price of the down-and-out call with the barrier at S = 0.9 as a function of the strike K; (b) a summary table. The number of MC paths for the MHP and MFD is 10,000; the number of paths for the MCS is 100,000.

minimum. Without loss of generality, we can restrict ourselves to the case of time-dependent drift and unit volatility by scaling time.⁵

To put things in perspective, we start with the standard Brownian motion and consider the following problem:

$$G_{t} - \frac{1}{2}G_{xx} = 0,$$

$$G(0, x) = \delta(x), \quad G(t, a) = 0,$$
(94)

where a is the lower bound, a < 0. It is clear that G(T, b; a) db is the probability of the Brownian motion ending in the interval (b - db/2, b + db/2) and having its minimum on the interval [a, 0]. Thus, the corresponding joint pdf has the form

$$\pi(T, a, b) = -\frac{\partial}{\partial a} G(T, b; a).$$
(95)

The method of images yields

$$G(T, b; a) = H(T, b) - H(T, b - 2a),$$
(96)

so that

$$\pi(T, a, b) = \frac{2}{T} (b - 2a) H(T, b - 2a).$$
(97)

For the drifted Brownian motion the problem can be written as follows:

$$G_{t} + \lambda G_{x} - \frac{1}{2}G_{xx} = 0,$$

$$G(0, x) = \delta(x), \quad G(t, a) = 0,$$
(98)

$$G(T,b;a) = H(T,b-\lambda T) - e^{2\lambda a}H(T,b-\lambda T-2a), \qquad (99)$$

$$\pi (T, a, b) = -\frac{\partial}{\partial a} G (T, b; a)$$

= $\frac{2}{T} (b - 2a) e^{2\lambda a} H (T, b - \lambda T - 2a).$ (100)

Needless to say that for zero drift, $\lambda = 0$, we recover Eq. (97).

Expressions (97) and (100) are very well-known, even though their derivation is usually somewhat convoluted. However, to the best of our knowledge, what is not known, is a similar expression when the drift λ depends on the (scaled) time v. We shall use the MHP to derive the corresponding formula. The problem of interest has the form

$$G_{\upsilon} + \lambda(\upsilon) G_{x} - \frac{1}{2}G_{xx} = 0, G(0, x) = \delta(x), \quad G(\upsilon, a) = 0,$$
(101)

$$G(\Upsilon, b; a) = H(\Upsilon, b - N_{\Upsilon}) - F(\Upsilon, b; a), \qquad (102)$$

$$F(\Upsilon, b; a) = \int_{0}^{\Upsilon} \frac{\left(b + \Theta(\Upsilon, \upsilon')(\Upsilon - \upsilon')\right) \exp\left(-\frac{\left(b + \Theta(\Upsilon, \upsilon')(\Upsilon - \upsilon')\right)^{2}}{2(\Upsilon - \upsilon')}\right)}{\sqrt{2\pi\left(\Upsilon - \upsilon'\right)^{3}}} \phi(\upsilon'; a) d\upsilon'$$
(103)

 $^{^{5}}$ Despite this being a classical problem, its correct solution is not presented in the literature, except for the simple case of constant drift. At the same time, several incorrect solutions have been proposed.

$$\phi(\upsilon; a) + \int_0^{\upsilon} \frac{\Theta(\upsilon, \upsilon') \Xi(\upsilon, \upsilon')}{\sqrt{2\pi(\upsilon - \upsilon')}} \phi(\upsilon'; a) \, d\upsilon' = f(\upsilon; a) \,, \tag{104}$$

$$f(v;a) = H(v, -N_v - a) = H(v, N_v + a).$$
(105)

Thus,

$$\pi\left(\Upsilon,a,b\right) = -\frac{\partial}{\partial a}G\left(\Upsilon,b;a\right)$$
$$= \int_{0}^{\Upsilon} \frac{\left(b + \Theta(\Upsilon,\upsilon')(\Upsilon-\upsilon')\right) \exp\left(-\frac{\left(b + \Theta(\Upsilon,\upsilon')(\Upsilon-\upsilon')\right)^{2}}{2(\Upsilon-\upsilon')}\right)}{\sqrt{2\pi(\Upsilon-\upsilon')^{3}}}\psi\left(\upsilon';a\right)d\upsilon',$$
(106)

where

$$\psi(\upsilon; a) + \int_0^{\upsilon} \frac{\Theta(\upsilon, \upsilon') \Xi(\upsilon, \upsilon')}{\sqrt{2\pi(\upsilon - \upsilon')}} \psi(\upsilon'; a) \, d\upsilon' = \frac{(N_\upsilon + a)}{\upsilon} f(\upsilon, a) \,. \tag{107}$$

6 Conclusions

This paper introduced a new hybrid MHP/MCS technique for pricing barrier options on assets with stochastic volatility. The idea is to decompose the solution process into the inner step, which solves a barrier problem for the conditionally independent process, and the outer step, which averages the corresponding solutions over the one-dimensional stochastic volatility dynamics.

Our methodology is general and can manage all known stochastic volatility models equally efficiently. Besides, relatively simple extensions (which will be described elsewhere) can also handle rough volatility models. With minimal changes, one can use the method to price popular double-no-touch options and other similar instruments.

While several authors used hybrid techniques before, see, e.g., Loeper and Pironneau (2009), their methods use the FDM and are still relatively slow, although undeniably faster than the standard two-dimensional MCS. Our method reduces the inner barrier problem to solving a linear Volterra equation of the second kind. It is very efficient and is an order of magnitude faster than other hybrid methods with the same (or better) accuracy. Our results are a natural generalization of Willard's formula, see Willard (1997), for barrier options.

As a byproduct of our analysis, we derived a new expression for the joint pdf for the value of a drifted Brownian motion and its running minimum or maximum in the case of time-dependent drift.

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